



Coupling of MD with Continuum Mechanics via a Bridging Scale Approach

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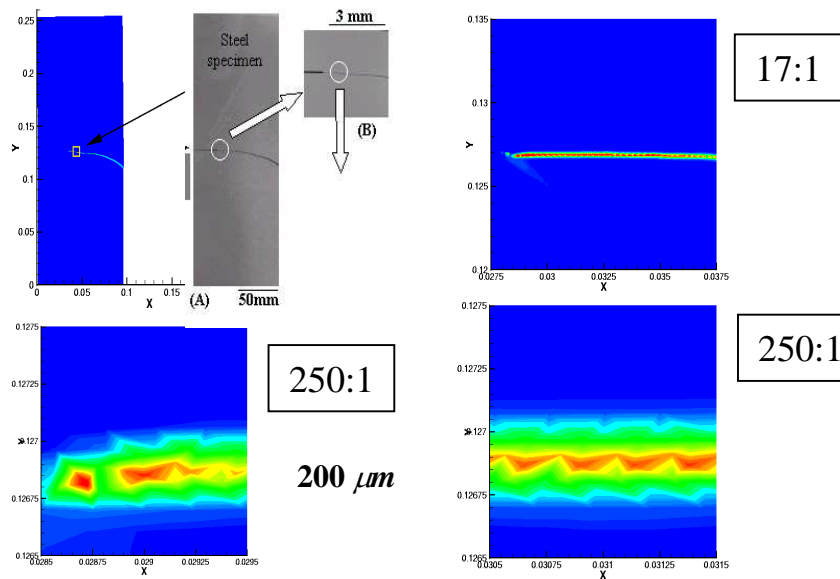
Collaborators

- Northwestern University
 - Prof. Wing Kam Liu
 - Prof. Dong Qian (U. of Cincinnati)
 - Harold Park
 - Eduard Karpov
 - Hiroshi Kadowaki
 - Sulin Zhang
 - Prof. Shaofan Li (U.C.-Berkeley)
- Sandia National Laboratories
 - Jonathan Zimmerman
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 - Patrick Klein



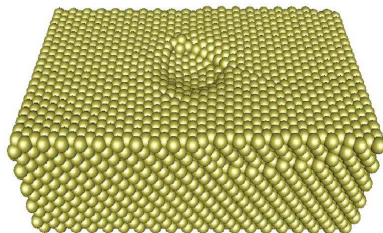
Examples of Multi-Scale Phenomena in Solids

Fracture/Failure of Solids



Prof. Shaofan Li, U.C.-Berkeley

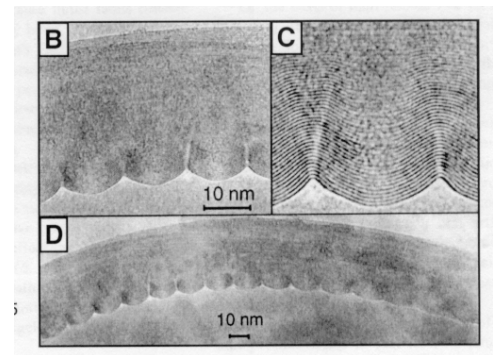
Nanoindentation



Dr. Eduard Karpov, Northwestern University

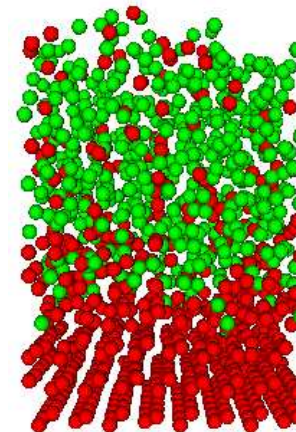
Nanoscale Devices

Poncharal et al., *Science* **283**:1513



Prof. Dong Qian, U. of Cincinnati

Film Growth





Concurrent Multiscales: Motivation

- Molecular dynamics simulations are limited to small domains ($\sim 10^6$ - 10^8 atoms) and small time frames (\sim nanoseconds)
 - Experiments, even on nano-systems, involve much larger systems over longer times
- Continuum models are good, but not always adequate
 - Problems in fracture and failure of solids require improved constitutive models to describe material behavior
 - Molecular dynamics is required in regions of high deformation or discontinuity
- Multiple scale nature of these problems calls for a combined molecular dynamics/continuum mechanics approach



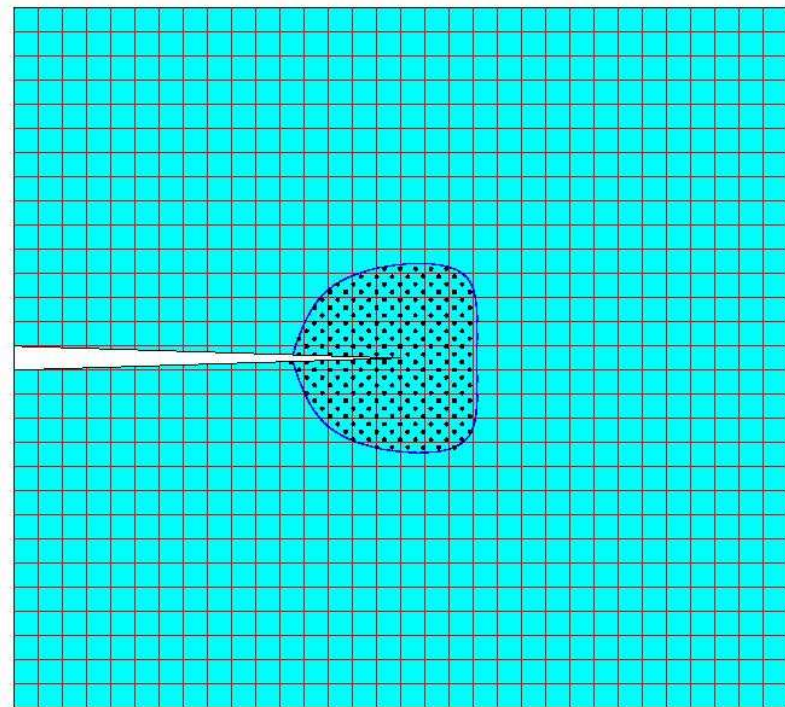
Concurrent Multiple Scales: Goals

- Method for coupling molecular dynamics to finite element or meshfree computations in concurrent simulations
 - Simulation of time dependent, finite temperature problems
- True “coarse scale” discretization in continuum
 - No meshing down to atomic scale
 - Subcycling time-stepping algorithms to take advantage of multiple time scales
 - don't want to be limited to nano time scale everywhere in the domain
- Easy implementation
 - Re-use of existing MD and continuum codes
 - Easily parallelizable algorithms



Concurrent Coupled Simulations

- Molecular dynamics to be used in region of interest
 - near crack/shear band tip
 - inside shear band
 - at area of large deformation
 - around dislocations
 - etc.
- Finite elements/meshless “coarse scale” defined *everywhere* in domain
 - not just overlap/handshake region
- Bridging scale used to ensure FEM gives correct coarse scale behavior





2-Part Strategy for Multiscale Coupling

- First, formally define exactly what is simulated at each scale
 - decompose total solution into coarse and fine scales
 - “bridging scale” used to represent the part of the total solution common to both simulations
 - provides coupling between the two simulations
- Second, eliminate fine scale degrees of freedom analytically outside of region of interest
 - use molecular dynamics (MD) only where necessary
 - use bridging scale decomposition to further define coupling between simulations
 - constitutive law in pure coarse scale region
 - boundary conditions on MD region



Coarse-Fine Decomposition

- Fields like displacement are decomposed into coarse and fine scales:

$$\mathbf{u} = \bar{\mathbf{u}} + \mathbf{u}'$$

- Coarse scale is represented by smooth basis functions associated with nodes, e.g. finite element shape functions:

$$\bar{\mathbf{u}}(\mathbf{X}_\alpha) = \sum_I N_I(\mathbf{X}_\alpha) \mathbf{d}_I \xrightarrow[\text{notation}]{\text{matrix}} \bar{\mathbf{u}} = \mathbf{N} \mathbf{d}$$

- nodal degrees of freedom minimize mass-weighted error norm:

$$E = (\mathbf{u} - \mathbf{N} \mathbf{d})^T \mathbf{M}_A (\mathbf{u} - \mathbf{N} \mathbf{d})$$

$$\frac{dE}{d\mathbf{d}} = 0 \Rightarrow \mathbf{d} = (\mathbf{N}^T \mathbf{M}_A \mathbf{N})^{-1} \mathbf{N}^T \mathbf{M}_A \mathbf{u}$$

- This leads to a definition of the coarse scale in terms of a projection matrix:

$$\boxed{\bar{\mathbf{u}} = \mathbf{P} \mathbf{u}}$$

$$\mathbf{P} = \mathbf{N} (\mathbf{N}^T \mathbf{M}_A \mathbf{N})^{-1} \mathbf{N}^T \mathbf{M}_A$$



Coarse-Fine Decomposition

- Once coarse scale is defined, fine scale is “everything else”:

$$\begin{aligned} \mathbf{u}' &= \mathbf{u} - \bar{\mathbf{u}} \\ &= \mathbf{u} - \mathbf{P}\mathbf{u} \\ &= (\mathbf{I} - \mathbf{P})\mathbf{u} \equiv \mathbf{Q}\mathbf{u} \end{aligned}$$

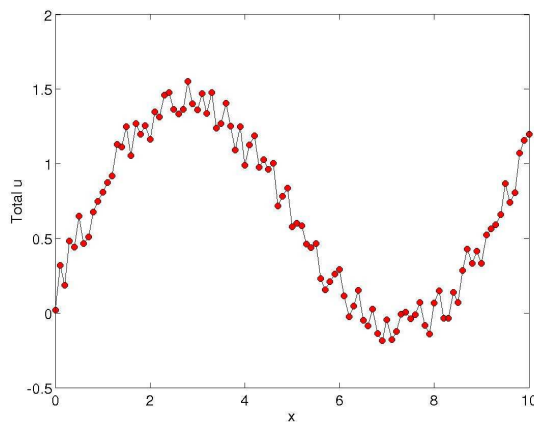
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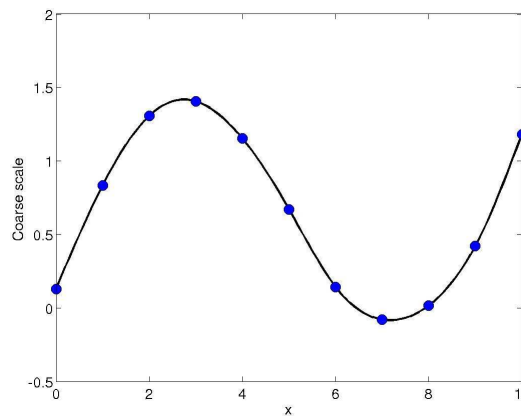
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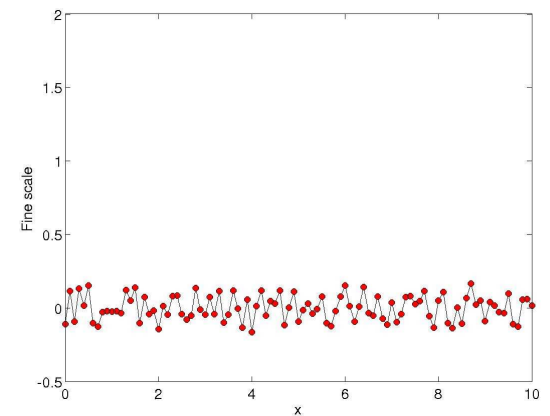
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Concurrent Multiscale Solution

- Use FEM for the coarse scale, MD for the fine scale in the decomposition:

$$\bar{\mathbf{u}} = \mathbf{N}\mathbf{d} \quad \longrightarrow \quad \text{solve } \mathbf{d} \text{ using FEM}$$

$$\mathbf{u}' = (\mathbf{I} - \mathbf{P})\mathbf{q} \quad \longrightarrow \quad \text{solve } \mathbf{q} \text{ using MD}$$

- With our choice of projection operator, kinetic energy separates completely into coarse and fine scales:

$$K_E = \frac{1}{2} \dot{\mathbf{d}}^T \mathbf{M} \dot{\mathbf{d}} + \frac{1}{2} \dot{\mathbf{q}}^T \tilde{\mathbf{M}} \dot{\mathbf{q}}$$

- Coupling between scales is only through the forcing term
 - Final momentum equations become:

$$\text{Coarse scale: } \mathbf{M} \ddot{\mathbf{d}} = \mathbf{N} \mathbf{f}(\bar{\mathbf{u}} + \mathbf{u}')$$

$$\text{Fine scale: } m_\alpha \ddot{\mathbf{q}}_\alpha = \mathbf{f}_\alpha(\bar{\mathbf{u}} + \mathbf{u}')$$

Reference: G. Wagner and W.K. Liu, *JCP* **190**:249-74 (2003).



Coarse Scale Modeling

- Coarse scale equation can be related to usual finite element treatment by approximating summations over atoms as domain integrals:

$$\sum_J \mathbf{M}_{IJ} \ddot{\mathbf{d}}_J = \mathbf{f}_I(\bar{\mathbf{u}}, \mathbf{u}')$$

where

$$\mathbf{M}_{IJ} = \int_{\Omega} \rho(\mathbf{x}) N_I(\mathbf{x}) N_J(\mathbf{x}) d\mathbf{x}$$

- The nodal force depends on the coarse scale only through the deformation gradient \mathbf{F} :

$$\begin{aligned} \mathbf{f}_I &= - \frac{\partial U}{\partial \mathbf{d}_I} \\ &= - \sum_{\alpha} \frac{\partial W_{\alpha}}{\partial \mathbf{d}_I} \Delta V_{\alpha} \\ &= - \sum_{\alpha} \frac{\partial \mathbf{F}_{\alpha}}{\partial \mathbf{d}_I} \frac{\partial W_{\alpha}}{\partial \mathbf{F}_{\alpha}^T} \Delta V_{\alpha} \approx - \int_{\Omega} N_{I,x}(\mathbf{x}) \mathbf{P}^K(\mathbf{x}) dV \end{aligned}$$



Fine Scale Boundary Conditions

- We want to avoid grading the coarse mesh down to the atomic lattice scale at the boundary
 - expensive
 - too much information
 - limits coarse scale time step
- Information passes from a fine MD lattice directly into a coarse scale mesh
 - small-scale energy can't be represented on the coarse scale, has nowhere else to go
 - leads to **internal reflection** of small scale waves
- Proper boundary treatment requires accounting for fine scale dynamics that are not simulated directly
 - correct boundary treatment falls out automatically from bridging scale decomposition
 - linearize in the fine scales at the boundary



MD Boundary Condition

$$\ddot{\mathbf{q}}_1(t) = \mathbf{M}_A^{-1} \mathbf{f}_1^*(t) + \int_0^t \boldsymbol{\theta}(t-\tau) \mathbf{a}'_2(\tau) d\tau + \mathbf{R}_1(t)$$
$$\mathbf{a}'_2(t) = \mathbf{M}_{A2}^{-1} \mathbf{f}_2^*(t) - \ddot{\mathbf{u}}_2(t)$$

Region 1: MD + FEM

Region 2: FEM only (+ “ghost atoms”)

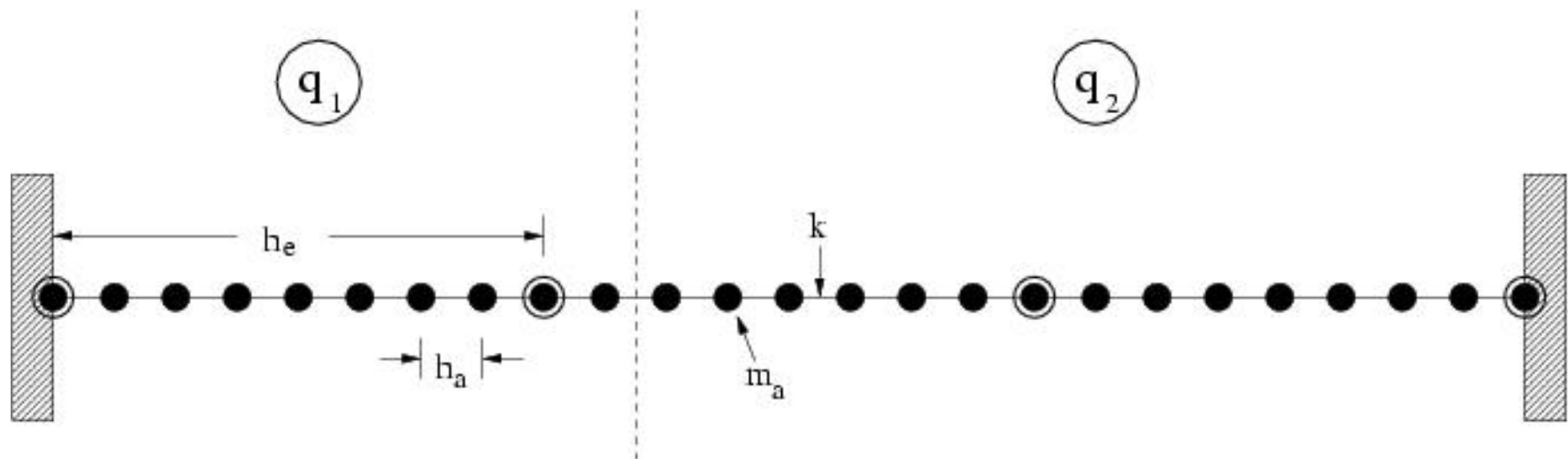
where $\mathbf{f}^*(t)$ are forces computed using just the coarse scale displacements outside MD region (e.g. through “ghost atoms”)

- The total forcing term consists of three major parts:
 - The standard interatomic force computed in MD simulation by assuming displacements of all atoms just outside the boundary are given by the coarse scale
 - A time history-dependent dissipation at the boundary (similar to a damping term)
 - A random forcing term at the boundary
 - the form of this term can be related to the temperature of the solid:

$$\langle R_i(t) R_j(0) \rangle = -\delta_{ij} \beta(t) k_B T$$



Example Problem: 1D Harmonic Chain

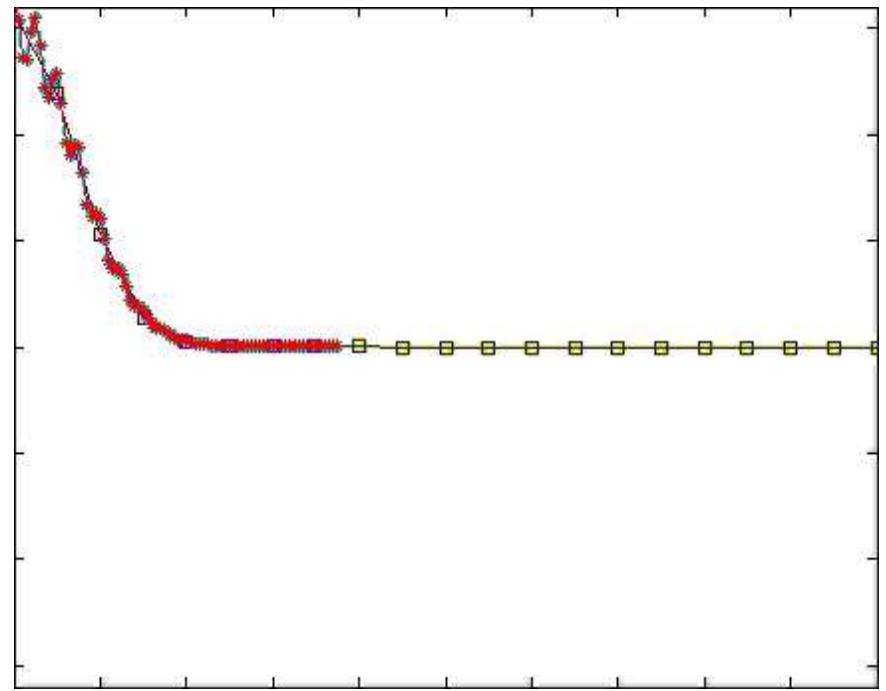
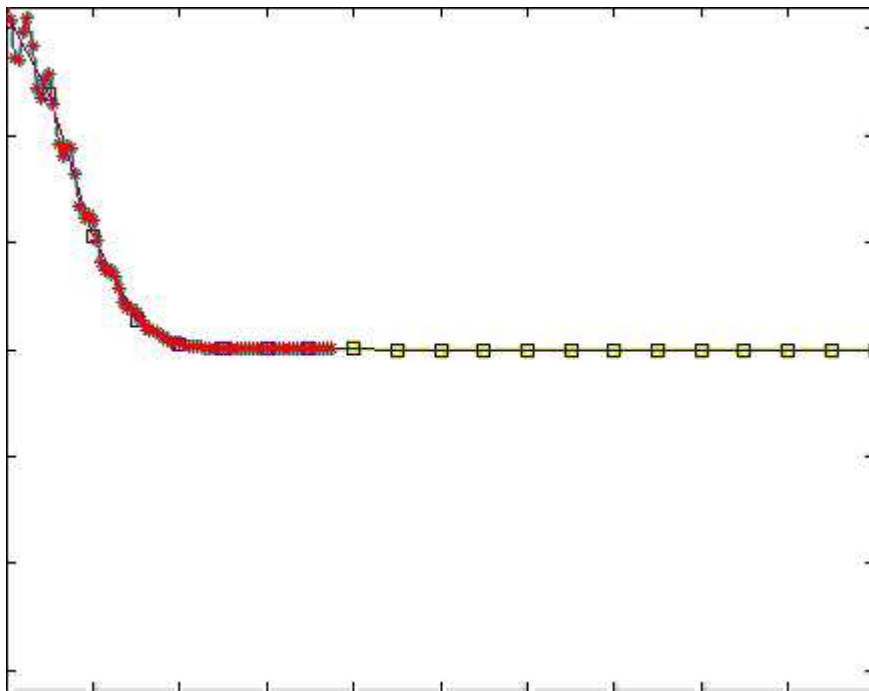




Effects of BC's on Internal Wave Reflection

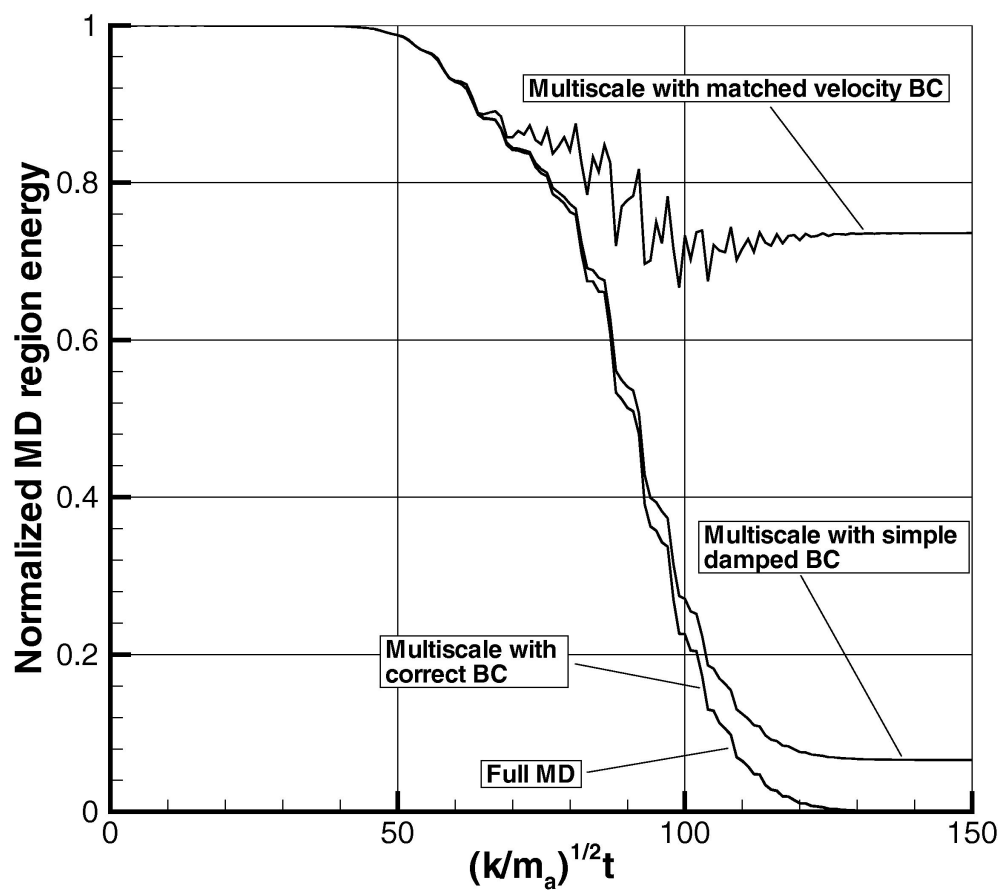
$$\dot{\mathbf{q}}_{\alpha} = \sum_I N_I(\mathbf{x}_{\alpha}) \dot{\mathbf{d}}_I$$

$$\ddot{\mathbf{q}}_1(t) = \mathbf{M}_A^{-1} \mathbf{f}_1^*(t) + \int_0^t \boldsymbol{\theta}(t-\tau) \mathbf{a}'_2(\tau) d\tau$$



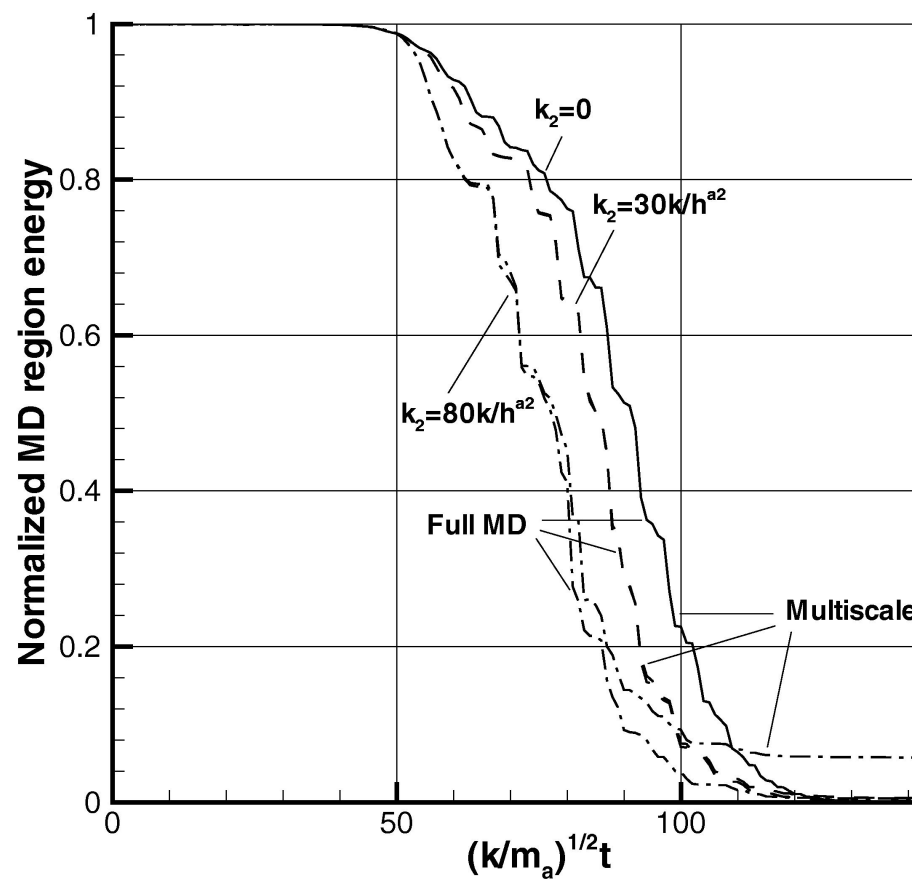


Energy Transfer out of MD Region





Energy Transfer out of MD Region: Nonlinear Potential



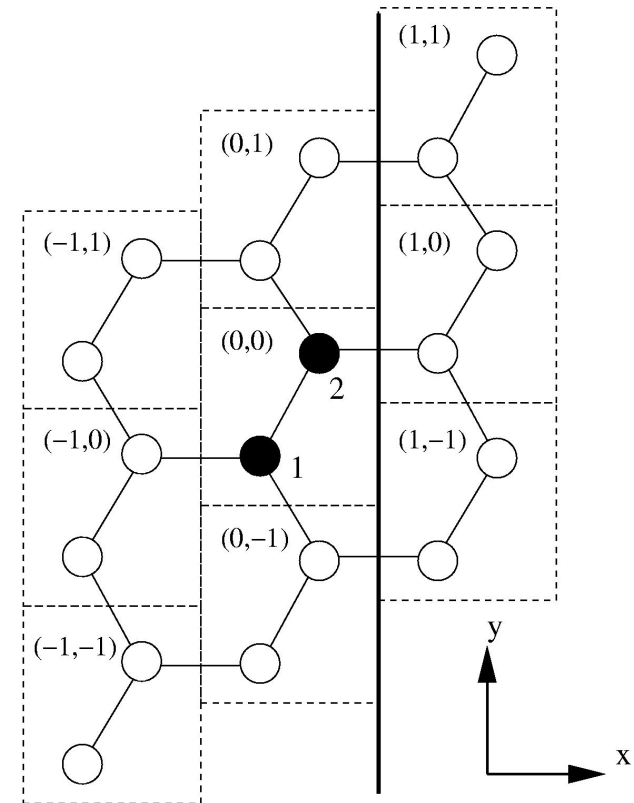


Damping Kernels in Multiple Dimensions

- Damping kernel can be easily computed for any regular crystal lattice across a planar boundary
 - periodicity allows spatial Fourier transform
 - unit cells can be indexed (l,m,n) in 3D, or (l,m) in 2D
 - boundary condition obtained by solving for atoms just outside boundary ($l=1$) in terms of atoms just inside boundary ($l=0$)
 - final boundary condition has form of **time history integration** with **spatial coupling** along boundary:

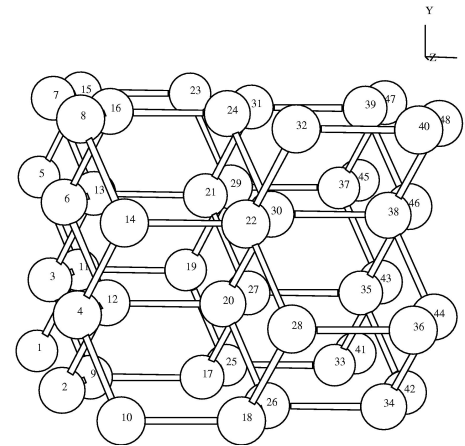
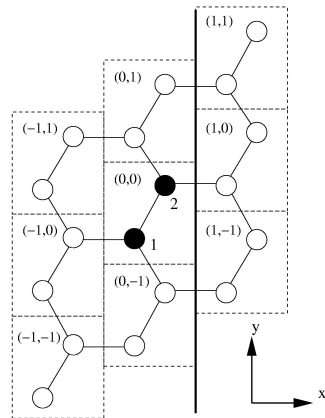
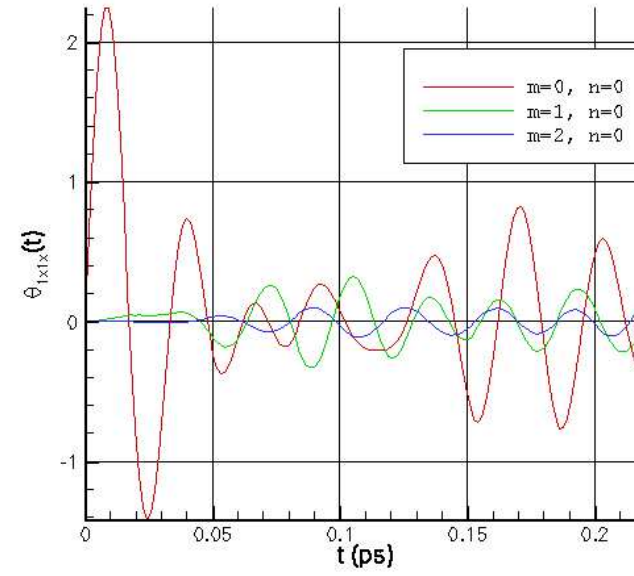
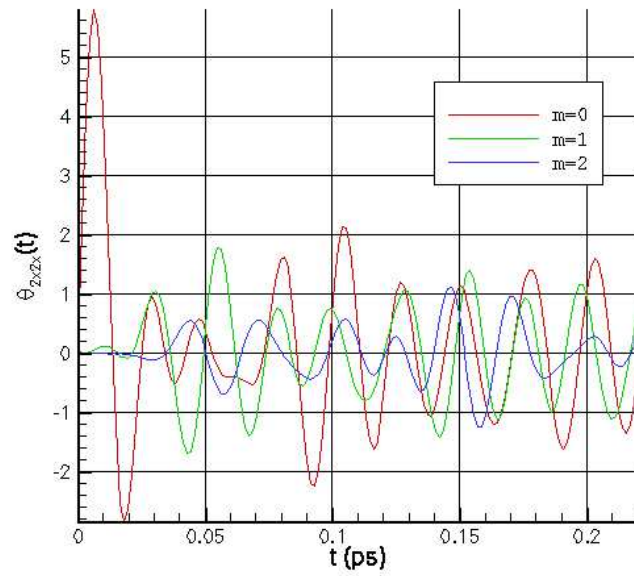
$$\mathbf{f}_m^{1 \rightarrow 0}(t) = \sum_{m'=-\infty}^{\infty} \int_0^t \boldsymbol{\theta}_{m-m'}(t-\tau) \mathbf{u}_{0,m'}(\tau) d\tau$$

- note that \mathbf{u} and \mathbf{f} are vectors containing all dof's in the unit cell, and $\boldsymbol{\theta}$ is a matrix coupling them
- Reference: G. Wagner, E. Karpov and W.K. Liu. *Comp Meth. Appl. Mech. Eng.*, to appear (2004).





Damping Kernels in 3D: Carbon Structures

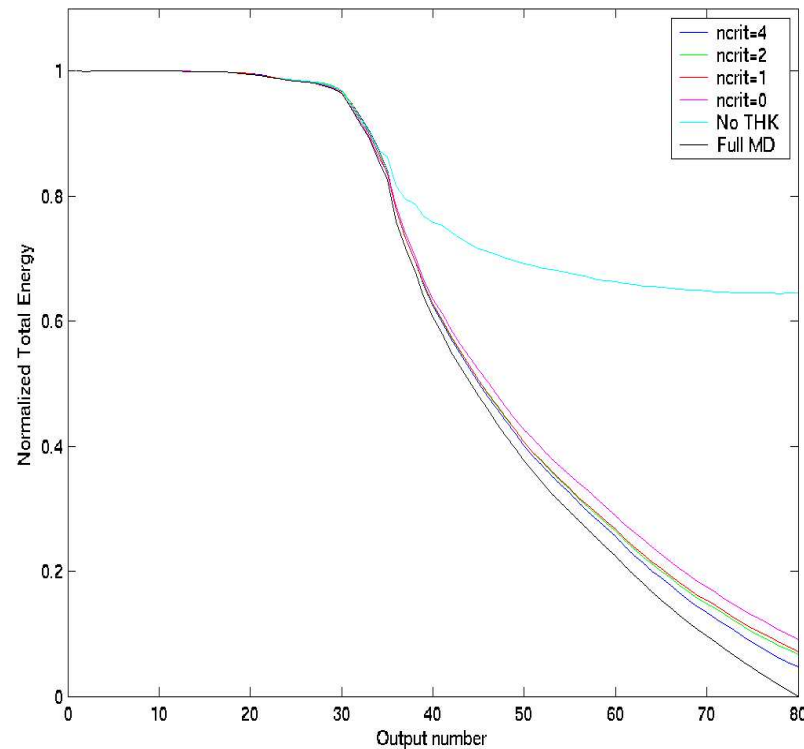




2D Multiscale Wave Propagation


smallnothk.mov


Smallthk.mov



Energy Transfer

Rates:

No BC: 35.47%

$N_{crit} = 0$: 90.94%

$N_{crit} = 1$: 92.85%

$N_{crit} = 2$: 93.34%

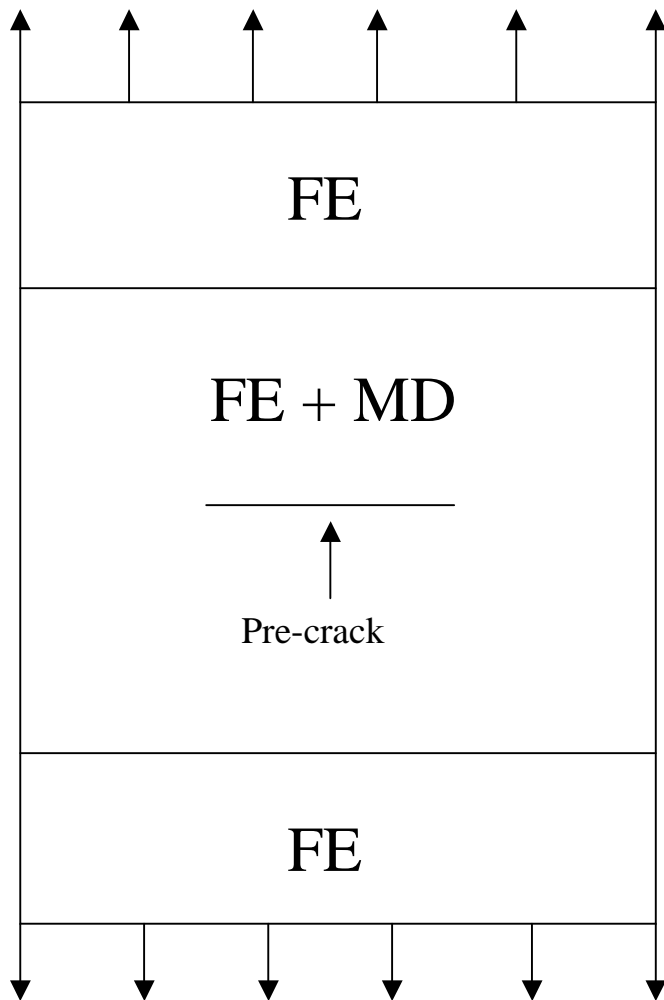
$N_{crit} = 4$: 95.27%

Full MD: 100%

Harold Park, Northwestern University



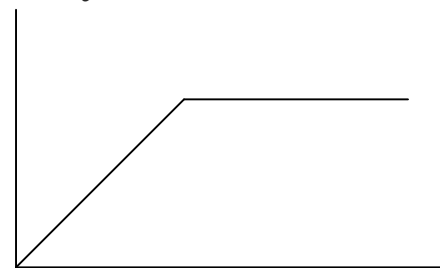
2D Dynamic Crack Propagation



Problem Description:

- LJ 6-12 potential, $\sigma=\epsilon=1$
- Nearest neighbor interactions
- 90000 atoms, 1800 finite elements (900 in coupled region)
 - 100 atoms per finite element
- $\Delta t_{fe} = 40\Delta t_{md}$
- Ramp velocity BC on FEM
- Full MD = 180,000 atoms

velocity



time



bsrampvel06_301x301_all.mov



Summary: Coupled MD/FEM

- Bridging scale decomposition allows concurrent simulation of fine scale using MD and coarse scale using FEM
 - bridging scale projection provides a unique decomposition of total solution for separation into coarse and fine scales
 - coarse scale mesh need not correspond to atomic lattice for coupling
 - subcycling can be used to take advantage of the different time scales in the coarse and fine regions
 - coarse scale equations and boundary conditions follow directly from the multi-scale formulation
- Future work:
 - study of approximations in boundary conditions
 - truncations of summations/integrals, approximations to kernel function
 - determination of most accurate/efficient integration of coarse scale region near MD boundary
 - development of coupled energy equation to track fine scale energy
 - time averaging as part of coarse scale projection